

## THE INTERSTELLAR 4.62 MICRON BAND

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### ABSTRACT

We present new 4.5–5.1  $\mu\text{m}$  (2210–1970  $\text{cm}^{-1}$ ) spectra of embedded protostars, W33 A, AFGL 961 E, AFGL 2136, NGC 7538 IRS 9, and Mon R2 IRS 2, which contain a broad absorption feature located near 4.62  $\mu\text{m}$  (2165  $\text{cm}^{-1}$ ), commonly referred to in the literature as the “X–C $\equiv$ N” band. The observed peak positions and widths of the interstellar band agree to within 2.5  $\text{cm}^{-1}$  and 5  $\text{cm}^{-1}$ , respectively. The strengths of the interstellar 4.62  $\mu\text{m}$  band and the ice absorption features in these spectra are not correlated, which suggests a diversity of environmental conditions for the ices we are observing. We explore several possible carriers of the interstellar band and review possible production pathways through far-ultraviolet photolysis (FUV), ion bombardment of interstellar ice analog mixtures, and acid-base reactions. Good fits to the interstellar spectra are obtained with an organic residue produced through ion bombardment of nitrogen-containing ices or with the OCN<sup>–</sup> ion produced either through acid-base reactions or FUV photolysis of NH<sub>3</sub>-containing ices.

*Subject headings:* infrared: ISM: lines and bands — ISM: clouds — ISM: molecules —  
line: identification — stars: pre-main-sequence

### 1. INTRODUCTION

More than two decades of infrared astronomical spectroscopy have revealed the importance of dust in the interstellar medium (ISM). Stars and planetary systems form out of dust in dense molecular clouds, the same dust that plays a pivotal role in the heating and cooling of the interstellar gas. Until relatively recently, crucial portions of the infrared spectrum were unavailable to astronomical spectroscopists owing to atmospheric limitations that hinder ground-based efforts, so despite the accumulating body of knowledge, gaps in the spectral observations have remained. The 2–30  $\mu\text{m}$  wavelength region is critically important for the study of organic solids and ices, and with results from the Kuiper Airborne Observatory (KAO) and the *Infrared Space Observatory* (ISO), the major gaps in the observational picture are now being filled in. The combination of moderate-resolution spectroscopy in those regions accessible from the ground with the space and airborne observations is providing a more complete inventory of the dust constituents, as well as information regarding the physical locations of dust repositories. Ultimately, the goal is to understand the physical conditions under which interstellar dust features form and evolve.

Here we present moderate-resolution ( $R \approx 1200$ ) infrared spectroscopy in a spectral region that is accessible from the ground. The feature discussed here, the interstellar 4.62  $\mu\text{m}$  (2165  $\text{cm}^{-1}$ ) band, likely arises from a carbon-nitrogen carrier and may be our most useful diagnostic of processed interstellar ices. The exact identification remains unclear, however, as there are multiple pathways to the production

of the band in the laboratory, and different nitrogen-containing precursor molecules have been used. Identifying the production mechanism responsible for the interstellar 4.62  $\mu\text{m}$  band and/or constraining the nitrogen reservoir available for processing will help our understanding of solid-state organic species in space.

The observation and reduction procedures used to obtain the astronomical spectra are discussed in § 2, and the results are presented in § 3. A discussion of possible identifications and formation mechanisms of the interstellar 4.62  $\mu\text{m}$  band are presented in § 4. A summary is given in § 5. The observational details are listed in Table 1, and a compilation of dust optical depth features is presented in Table 2. Optical depth ratios of diagnostic absorption features are provided in Table 3. Summary tables of the candidate species and production pathways are given in Tables 4 and 5, respectively.

### 2. OBSERVATIONS AND DATA REDUCTION

All the observations presented here were made during 1986 November 17–21 and 1988 July 21–22 at the NASA Infrared Telescope Facility (IRTF) on Mauna Kea using the Cooled Grating Array Spectrometer (CGAS). The CGAS employed a linear 32 element InSb detector (Tokunaga, Smith, & Irvin 1987). Output from detectors 1, 2, 13, and 32 was not used, as these detectors were not operating reliably during the observation run. The observations were made using an aperture of 2".7 and a north-south chop of 15'. Typical exposure times were 30–60 minutes per source per night. The spectrometer was used in first order with grating B, which provided a resolving power of 1200 (i.e., a spectral resolution element of 0.004  $\mu\text{m}$  [ $\approx 1.8 \text{ cm}^{-1}$ ]). Objects were observed at two grating settings separated by half a resolution element, except for NGC 7538 IRS 9, which was observed once (1989 July) with 0.25 resolution element sampling. For calibration and atmospheric correction purposes, each individual grating setting on an object was either preceded or followed by a standard star located in the same part of the sky that also had a closely matching air mass to that of the source. This process minimized the residual telluric absorption features

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TABLE 1  
SUMMARY OF OBSERVATIONS

Source	R.A. (1950)	Decl. (1950)	Date (UT)	Standard
AFGL 2136 .....	18 24 14	-25 06 12	1988 Jul 22	BS 7020
AFGL 961 E .....	6 31 59.1	4 15 10	1986 Nov 17-19	BS 1903
Mon R2 IRS 2 .....	06 05 19.4	-06 22 24	1986 Nov 17-19	BS 1903
W33 A .....	18 11 43.7	-17 53 02	1988 Jul 21	BS 6561
NGC 7538 IRS 9 .....	23 11 52.8	61 10 58	1986 Nov 17-19 and 1989 July 20	BS 8571

NOTE.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.

in the reduced spectra. After ratioed spectra were calculated at individual grating settings, the flux levels were adjusted to each other in the overlap region. Mismatches in flux levels, caused by seeing and guiding errors, were typically less than 10%.

The wavelength calibration is discussed in detail in Tielens et al. (1991). The central wavelength of each detector is known to much better precision than the spectral resolution element. Absolute wavelength calibration was provided by an argon discharge lamp attached to the CGAS. Measurements at the telescope show that wavelengths are accurate to about  $5 \times 10^{-4} \mu\text{m}$  ( $0.23 \text{ cm}^{-1}$ ), and therefore the shape of an absorption feature is determined to this accuracy. We note that two features separated by less than the Nyquist frequency ( $\approx 0.9 \text{ cm}^{-1}$ ) will not be resolved. All wavelengths are given here in vacuo. Doppler shifts

due to the motion of the object are typically very small, and no correction has been applied to the reported spectra.

### 3. RESULTS

#### 3.1. The Objects and Their Spectra

We present new observations of AFGL 2136, AFGL 961 E, Mon R2 IRS 2, W33 A, and NGC 7538 IRS 9. Table 1 contains a list of the objects and their coordinates. These sources are luminous, deeply embedded protostars seen through dense molecular cloud material. The objects in the present paper were chosen to sample the environments of different molecular cloud cores where other absorption features had previously been detected (specifically,  $\text{H}_2\text{O}$ ,  $\text{CO}$ , and silicate). The  $4.5\text{--}5.1 \mu\text{m}$  ( $2222\text{--}1960 \text{ cm}^{-1}$ ) spectrum of W33 A is shown in Figure 1. The  $4.5\text{--}4.8 \mu\text{m}$  ( $2222\text{--}2083$

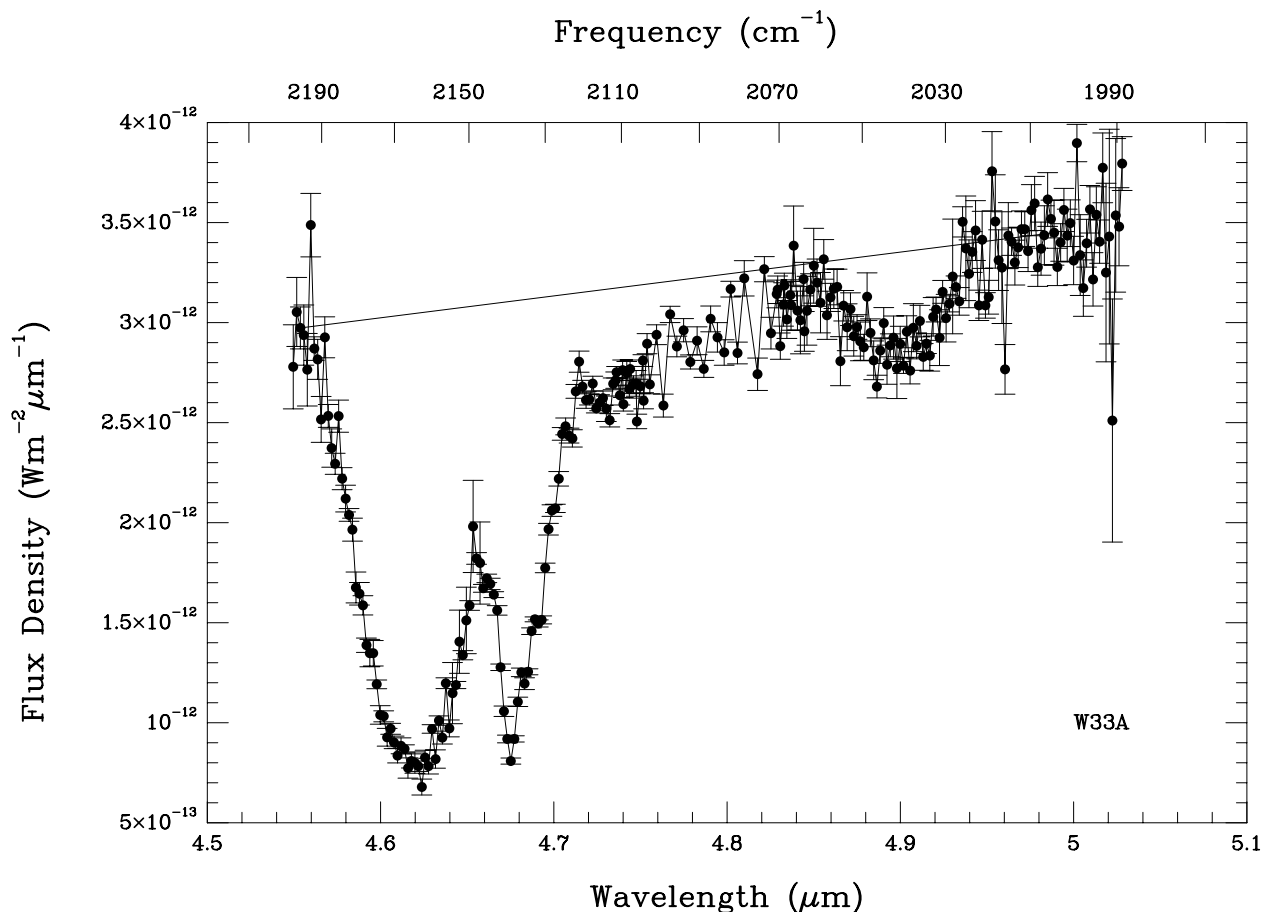


FIG. 1.—Flux spectra ( $\text{W m}^{-2} \mu\text{m}^{-1}$ ) of W33 A  $4.5\text{--}5.1 \mu\text{m}$  ( $2222\text{--}1960 \text{ cm}^{-1}$ ). The solid line depicts the continuum baseline from which the optical depth determinations were made.

TABLE 2  
OPTICAL DEPTH OF ABSORPTION BANDS

Source	$\tau_{3.1}$ (H <sub>2</sub> O)	$\tau_{3.47}$	$\tau_{3.5}$ Wing	$\tau_{4.62}$	$\tau_{6.8}$	$\tau_{4.67(\text{CO})}$ Nonpolar 2140 cm <sup>-1</sup>	$\tau_{4.67(\text{CO})}$ Polar 2136 cm <sup>-1</sup>	$\tau_{4.9}$ (OCS)	$\tau_{9.7}$ (Silicate)	Object $A_V$	$N(\text{H})$ <sup>a,b</sup> (cm <sup>-2</sup> )
AFGL 2136.....	2.72 <sup>e</sup>	...	...	0.06	...	0 <sup>b</sup>	0.22 <sup>b</sup>	...	5.07 <sup>c</sup>	94.3 <sup>d</sup>	$1.8 \times 10^{23}$
AFGL 961 E.....	2.46 <sup>e</sup>	0.68 <sup>f</sup>	0.36 <sup>e</sup>	0.06	...	0.22 <sup>g</sup>	0.20 <sup>g</sup>	<0.02	2.11 <sup>d</sup>	39 <sup>b</sup>	$7.4 \times 10^{22}$
W33 A.....	5.4 <sup>d</sup>	0.15 <sup>h</sup>	2.53 <sup>d</sup>	1.5	1.8 <sup>i</sup>	0.79 <sup>g</sup>	0.68 <sup>g</sup>	0.14 <sup>j</sup>	7.84 <sup>d</sup>	145 <sup>k</sup>	$28 \times 10^{22}$
Elias 18.....	0.80 <sup>l</sup>	0.022 <sup>f</sup>	...	0.05 <sup>m</sup>	...	0.26 <sup>g</sup>	0.14 <sup>g</sup>	...	0.43 <sup>l</sup>	17-20 <sup>k</sup>	$2.2 \times 10^{17}$
Mon R2 IRS 2.....	2.54 <sup>e</sup>	0.046 <sup>f</sup>	0.46 <sup>e</sup>	0.029	...	0.37 <sup>g</sup>	0.16 <sup>g</sup>	0.015 <sup>n</sup>	2.11 <sup>b</sup>	21 <sup>b</sup>	$4 \times 10^{22}$
L1551 IRS 5.....	2.1 <sup>o</sup>	...	...	0.16 <sup>m</sup>	...	0.80 <sup>m</sup>	0.16 <sup>m</sup>	...	1.4 <sup>b</sup>	12.1 <sup>p</sup>	$2.3 \times 10^{22}$
Elias 16.....	1.6 <sup>m</sup>	<0.031 <sup>q</sup>	<0.016 <sup>c</sup>	<0.05 <sup>m</sup>	...	1.22 <sup>r</sup>	0.11 <sup>g</sup>	<0.04	0.66 <sup>l</sup>	22.3 <sup>k</sup>	$3.8 \times 10^{20}$
NGC 7538 IRS 9.....	3.28 <sup>d</sup>	0.10 <sup>f</sup>	0.64 <sup>d</sup>	0.31	0.3 <sup>r</sup>	2.50 <sup>g</sup>	0.2 <sup>g</sup>	...	4.46 <sup>d</sup>	84 <sup>s</sup>	$16 \times 10^{22}$

<sup>a</sup> From  $N(\text{H})/A_V = 1.9 \times 10^{21}$ .

<sup>b</sup> Tielens et al. 1991.

<sup>c</sup> Hough et al. 1989.

<sup>d</sup> Willner et al. 1982.

<sup>e</sup> Smith, Sellgren, & Tokunaga 1989.

<sup>f</sup> Brooke, Sellgren, & Smith 1996.

<sup>g</sup> Chiar et al. 1998.

<sup>h</sup> Allamandola et al. 1992.

<sup>i</sup> Tielens & Allamandola 1987.

<sup>j</sup> Palumbo, Tielens, & Tokunaga 1995.

<sup>k</sup> Averaged from Elias 1978 and Smith, Sellgren, & Brooke 1993.

<sup>l</sup> Whittet et al. 1988.

<sup>m</sup> Tegler et al. 1995.

<sup>n</sup> Palumbo et al. 1998.

<sup>o</sup> Sato et al. 1990.

<sup>p</sup> Elias 1978.

<sup>q</sup> Chiar, Adamson, & Whittet 1996.

<sup>r</sup> Schutte et al. 1996.

<sup>s</sup> Visual extinction in units of magnitudes from  $\tau_{s_i} = A_V/18.6$ .

TABLE 3  
 OPTICAL DEPTH RATIOS

Source	$\tau_{4.67(\text{CO})}$ Nonpolar/Polar	$\tau_{4.62}/\tau_{4.67(\text{CO})}$ Nonpolar	$\tau_{4.62}/\tau_{4.67(\text{CO})}$ Polar	$\tau_{4.62}/A_V$	$\tau_{4.62}/\tau_{3.1}$	$\tau_{4.62}/\tau_{9.75}$	$\tau_{9.7}/\tau_{3.1}$
AFGL 2136 .....	0	0	0.27	0.001	0.022	0.012	1.86
AFGL 961 E .....	1.1	0.27	0.30	0.002	0.024	0.028	0.86
W33A .....	1.2	1.9	2.21	0.01	0.278	0.191	1.45
Elias 18 .....	1.9	0.19	0.36	0.003	0.063	0.116	0.54
Mon R2 IRS 2 .....	2.3	0.078	0.18	0.001	0.011	0.014	0.83
L1551 IRS 5 .....	5.0	0.20	1.0	0.013	0.076	0.114	0.67
Elias 16 .....	11.1	0.041	0.45	0.002	0.031	0.076	0.41
NGC 7538 IRS 9 .....	12.5	0.124	1.55	0.004	0.095	0.070	1.36

$\text{cm}^{-1}$ ) spectra of AFGL 961, AFGL 2136, AFGL 2591, Mon R2 IRS 2, and NGC 7538 IRS 9 are shown in Figures 2 and 3. The spectra show features at  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) and  $4.67 \mu\text{m}$  ( $2140 \text{ cm}^{-1}$ ). In the case of W33 A, the broader spectral coverage reveals another absorption feature at  $4.9 \mu\text{m}$  ( $2041 \text{ cm}^{-1}$ ). The  $4.67 \mu\text{m}$  and  $4.9 \mu\text{m}$  features are attributed to absorption by solid-phase carbon monoxide and carbonyl sulfide, and these are discussed elsewhere (Tielens et al. 1991 and Palumbo, Tielens, & Tokunaga 1995, respectively). A recent discussion of CO in these sources can be found in Chiar et al. (1998).

For the  $4.62 \mu\text{m}$  and  $4.67 \mu\text{m}$  features, a straight-line continuum was constructed between the data points near  $4.56$  and  $4.76 \mu\text{m}$  ( $2190$  and  $2100 \text{ cm}^{-1}$ ). In the case of W33 A, where the observational coverage extended to  $5.1 \mu\text{m}$ , the continuum was drawn from  $4.55$  to  $4.98 \mu\text{m}$  ( $2197$ – $2008 \text{ cm}^{-1}$ ). The optical depths derived from the W33 A spectra are shown in Figure 4, and optical depth plots for AFGL 961, AFGL 2136, NGC 7538 IRS 9, and Mon R2 IRS 2 are

presented in Figures 5 and 6. Peak optical depths are listed in Table 2.

The profiles of the  $4.62 \mu\text{m}$  band in W33 A and AFGL 961 are compared in Figure 7a. These are the only two sources with a strong enough  $4.62 \mu\text{m}$  feature for such a comparison to be meaningful. Due to the strong solid CO band, there is an insufficient continuum to apply a  $\chi^2$  Gaussian fitting routine to these profiles. However, the peak positions in these sources agree to within  $2.5 \text{ cm}^{-1}$ , and the width of the feature agrees to within  $5 \text{ cm}^{-1}$ .

Gas-phase CO lines are known to be present in the spectra of many embedded objects (see Tielens et al. 1991 for discussion). Figure 3a shows the spectrum of the source AFGL 2591, in which the gas-phase CO lines are readily identified. A comparison to the spectra of Mon R2 IRS 2 and NGC 7538 IRS 9 (Fig. 3b and 3c, respectively), demonstrates the presence of weak CO lines in these sources as well. At low spectral resolution, blending of the rovibrational lines can give rise to broad absorption features that may

 TABLE 4  
 SUMMARY OF CANDIDATE SPECIES

Candidate Molecule	Comments	Feasibility
Nitriles ( $-\text{C}\equiv\text{N}$ ):		
Dicyandiamide [ $\text{NHC}(\text{NH}_2)\text{NHCN}$ ] .....	Band at $2165 \text{ cm}^{-1}$ but stronger $2204 \text{ cm}^{-1}$ band is not observed.	Poor
Acetonitrile (methylocyanide) $\text{CH}_3\text{CN}$ .....	Band at $2264 \text{ cm}^{-1}$ in an $\text{H}_2\text{O}$ matrix.	Poor
Pyruvonnitrile ( $\text{CH}_3\text{COCN}$ ) .....	Band at $2092 \text{ cm}^{-1}$ in $\text{H}_2\text{O}$ and $2225 \text{ cm}^{-1}$ in argon.	Poor
Cyanogen ( $\text{C}_2\text{N}_2$ ) .....	Band absorbs at $2167 \text{ cm}^{-1}$ in solid state, but lab spectra show a band at $240 \text{ cm}^{-1}$ 2 orders of magnitude stronger than $2165 \text{ cm}^{-1}$ ; not seen in ISM.	Poor
Thiocyanates ( $\text{S}-\text{C}\equiv\text{N}$ ) .....	No studies available.	Unknown
Isonitriles ( $-\text{N}\equiv\text{C}$ ):		
Methylisocyanide (isoacetonitrile) ( $\text{CH}_3\text{NC}$ ) .....	Typically stronger than nitriles, but less stable; might also expect the corresponding nitrile to be present. Absorbs at $2161 \text{ cm}^{-1}$ in both gas phase and argon matrix; might expect related $\text{CH}_3\text{CN}$ band at $2264 \text{ cm}^{-1}$ ; not seen in ISM.	Weak
<i>N</i> -butylisocyanide [ $\text{CH}_3(\text{CH}_2)_3\text{NC}$ ] .....	Strong CN stretch near $2165 \text{ cm}^{-1}$ in $\text{H}_2\text{O}$ and argon matrices.	Weak
Methylisocynoacetate ( $\text{CH}_3\text{OCOCH}_2\text{NC}$ ) .....	Strong CN stretch near $2165 \text{ cm}^{-1}$ in $\text{H}_2\text{O}$ and argon matrices.	Weak
Isothiocyanates ( $\text{S}-\text{N}\equiv\text{C}$ ) .....	No studies of isothiocyanates available.	Unknown
Cumulative double bonds ( $\text{C}=\text{C}=\text{O}$ ):		
Ketene $\text{H}_2\text{C}_2\text{O}$ .....	Absorbs at $2142 \text{ cm}^{-1}$ in argon or $\text{O}_2$ ; need to see behavior in a more appropriate matrix.	Unknown
Cyanate ions ( $\text{OCN}^-$ ) .....	Band falls in $2165$ – $2222 \text{ cm}^{-1}$ region; position sensitive to specific cation and solvation effects.	Good
Silanes:		
Organic silane ( $-\text{SiH}_3$ ) .....	$\text{Si}-\text{H}$ stretch falls in the $2160$ – $2140 \text{ cm}^{-1}$ range. Position varies with molecular group; possible that one can be found to match ISM feature.	Weak
Silane ( $\text{SiH}_4$ ) .....	Bands at $2216$ , $2176 \text{ cm}^{-1}$ in argon and $2200$ , $2165 \text{ cm}^{-1}$ in $\text{H}_2\text{O}$ ; only one band detected in the ISM.	Poor

TABLE 5

## PRODUCTION PATHWAYS

Starting Molecule	Product	Comments	References
UV Photolysis of Ices			
CO, NH <sub>3</sub> H <sub>2</sub> O, CH <sub>3</sub> OH ices .....	OCN <sup>-</sup>	Presence of ions supported by experiments using electron donors/acceptors; C,N,O are implicated in the carrier	1, 2, 3, 4, 5, 6, 7, 8, 9, 10
N <sub>2</sub> mixtures .....	No 2165 cm <sup>-1</sup> band		11
Grain Surface Reactions			
Hydrogenation/nitrogen addition:			
H + CO .....	HCO		12
N + HCO .....	NCHO* to HNCO	Stabilization of the excited NCHO* to its lowest energy level, HNCO	13, 14
Acid-base chemistry:			
H <sub>2</sub> O, NH <sub>3</sub> , HNCO .....	OCN <sup>-</sup>	Without UV. Deposition onto 12 K substrate yields OCN <sup>-</sup> and NH <sub>4</sub> <sup>+</sup> , band at 2171 cm <sup>-1</sup> grows during warm-up.	15
Ion Bombardment of Ices			
H <sub>2</sub> O:CH <sub>4</sub> :NH <sub>3</sub> or N <sub>2</sub> ices (2:1:1) .....	R-O-CN	30-60 keV helium, argon ions; produces 2165 cm <sup>-1</sup> band. R is a species belonging to the refractory residue	16
H <sub>2</sub> O:NH <sub>3</sub> :CH <sub>4</sub> .....		(≈1:3:2); MeV protons; produces 2170 cm <sup>-1</sup> band	17
H <sub>2</sub> O:N <sub>2</sub> :CO <sub>2</sub> .....		(≈5:1:1 and 1:1:1); band at 2180 cm <sup>-1</sup>	17

REFERENCES.—(1) Hagen 1982; (2) d'Hendecourt et al. 1986; (3) Allamandola et al. 1988; (4) Sandford & Allamandola 1990; (5) Bernstein et al. 1997; (6) Grim 1988; (7) Demyk et al. 1999; (8) Grim & Greenberg 1987b, (9) Grim et al. 1989a; (10) Schutte & Greenberg 1997; (11) Eisila et al. 1997; (12) van IJzendoorn et al. 1983; (13) Jacox & Milligan 1964; (14) Bernstein 1998; (15) Keane & Schutte 1998; (16) Palumbo et al. 1998; (17) Moore et al. 1983.

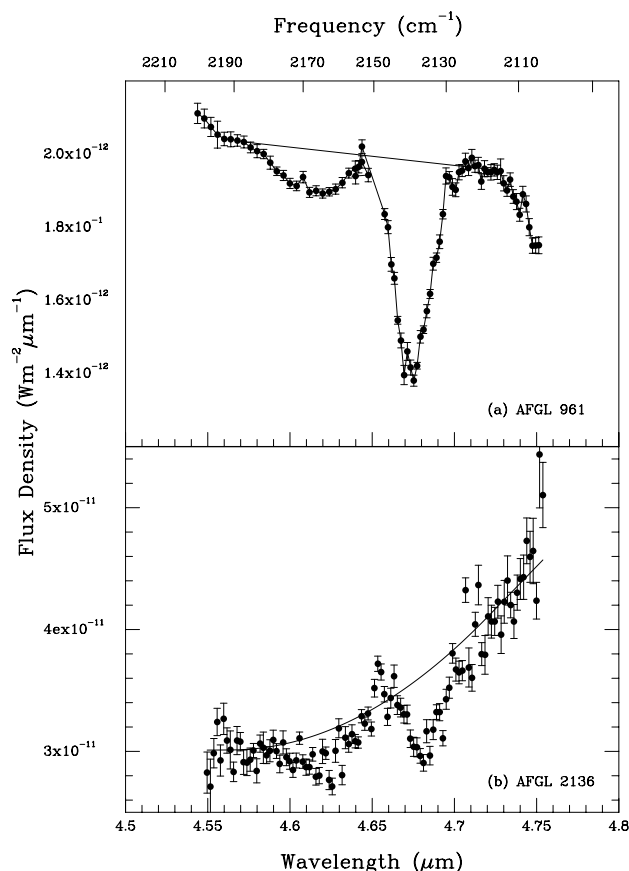


FIG. 2.—Flux spectra ( $\text{W m}^{-2} \mu\text{m}^{-1}$ ) of (a) AFGL 961 and (b) AFGL 2136 from 4.5–4.8  $\mu\text{m}$  (2222–2083  $\text{cm}^{-1}$ ). The solid line depicts the continuum baseline from which the optical depth determinations were made.

mimic solid-state absorptions (see Geballe 1986). However, at our resolution, the individual rovibrational lines are separated by more than one resolution element (see Fig. 3a), and blending will only occur for unreasonably large broadening parameters, i.e.,  $\delta v > 500 \text{ km s}^{-1}$ . Indeed, this is well in excess (by factors  $\approx 100$ ) of the width of CO individual rovibrational lines measured toward W33 A, NGC 7538 IRS 9, AFGL 2136, and AFGL 961 E (Mitchell et al. 1990). Moreover, there is no sign in our spectra of a corresponding *P*-branch of CO on the longer wavelength side of the 4.62  $\mu\text{m}$  band, which should be equally strong compared to the *R*-branch if the 4.62  $\mu\text{m}$  band were a result of gas-phase CO.

### 3.2. Interrelationships of Absorption Features

As shown in Table 2, the strength of the 4.62  $\mu\text{m}$  band per unit  $A_V$  (total visual extinction) varies by about an order of magnitude from cloud to cloud. Apparent from Table 3 is the fact that the optical depth of the 4.62  $\mu\text{m}$  band does not correlate with any other feature. The sources in Tables 2 and 3 are ordered from top to bottom in terms of decreasing temperature, as deduced from the amount of the volatile CO component (i.e., a nonpolar matrix) present relative to the amount of polar CO. The lack of correlation of the 4.62  $\mu\text{m}$  band to the ice features may be indicative of (1) the location of the carrier of this material (i.e., circumstellar rather than intracloud dust), (2) the degree of processing to which the dust has been exposed, and/or (3) differences in the evolutionary stages of the embedded sources. It should

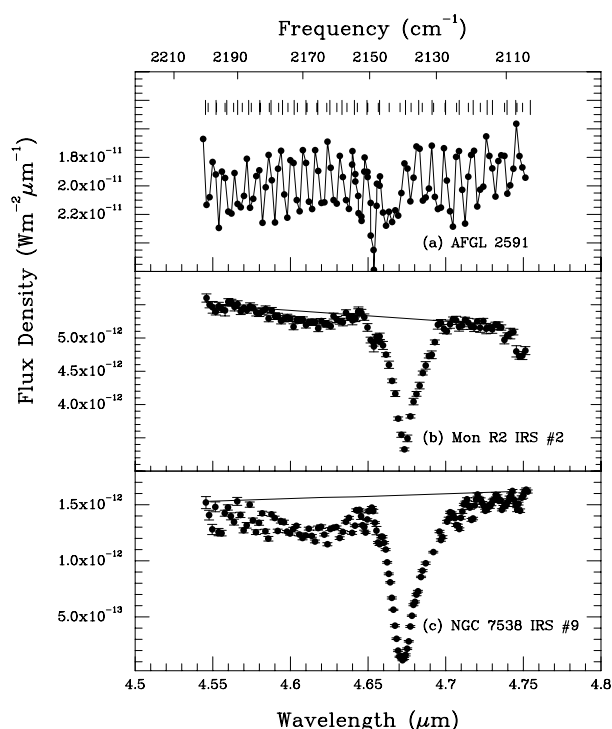


FIG. 3.—Flux spectra ( $\text{W m}^{-2} \mu\text{m}^{-1}$ ) of (a) AFGL 2591, (b) Mon R2 IRS 2, and (c) NGC 7538 IRS 9 from 4.5–4.8  $\mu\text{m}$  (2222–2083  $\text{cm}^{-1}$ ). Gas-phase CO lines ( $^{12}\text{CO}$  and  $^{13}\text{CO}$ ) have previously been identified in AFGL 2591 (Geballe & Wade 1985; Tielens et al. 1991) and are shown here for comparison to NGC 7538 IRS 9 and Mon R2 IRS 2 to illustrate the presence of CO gas-phase lines in the spectra of those sources. The solid line depicts the continuum baseline from which the optical depth determinations were made.

be noted that none of the other features appear to correlate with each other either, which may simply be a reflection of the fact that the formation of these features is much more critically dependent upon the environmental conditions than was previously thought.

### 4. POSSIBLE IDENTIFICATIONS OF THE 4.62 MICRON BAND

The 4.62  $\mu\text{m}$  band (Soifer et al. 1979; Lacy et al. 1984; Tegler et al. 1993, 1995), is commonly referred to in the literature as the X–C $\equiv$ N band. The association of the interstellar feature with nitriles (–C $\equiv$ N) came from early identifications with laboratory spectra thought to contain nitrile or isonitrile (–N $\equiv$ C) groups. Although X– commonly refers to a halogen in chemistry circles, the use of X– in the astronomical literature simply denotes some undetermined species; we adopt this nomenclature here as well. The interstellar band has also been attributed to the cyanate ion (OCN $^-$ ) by several laboratory groups (Grim & Greenberg 1987a; Grim et al. 1989b; Schutte & Greenberg 1997; Demyk et al. 1999). Other candidates have been suggested in the literature as well, and they are summarized in Table 4. In some cases the proposed species has a feature that does not coincide with the position of the interstellar band, but the species has not been tested in an astrophysically relevant matrix. Until those experiments have been conducted, the candidate cannot be assessed completely.

Triple bond compounds including nitriles (X–C $\equiv$ N), isonitriles (X–N $\equiv$ C), thiocyanates (X–S–C $\equiv$ N), iso-

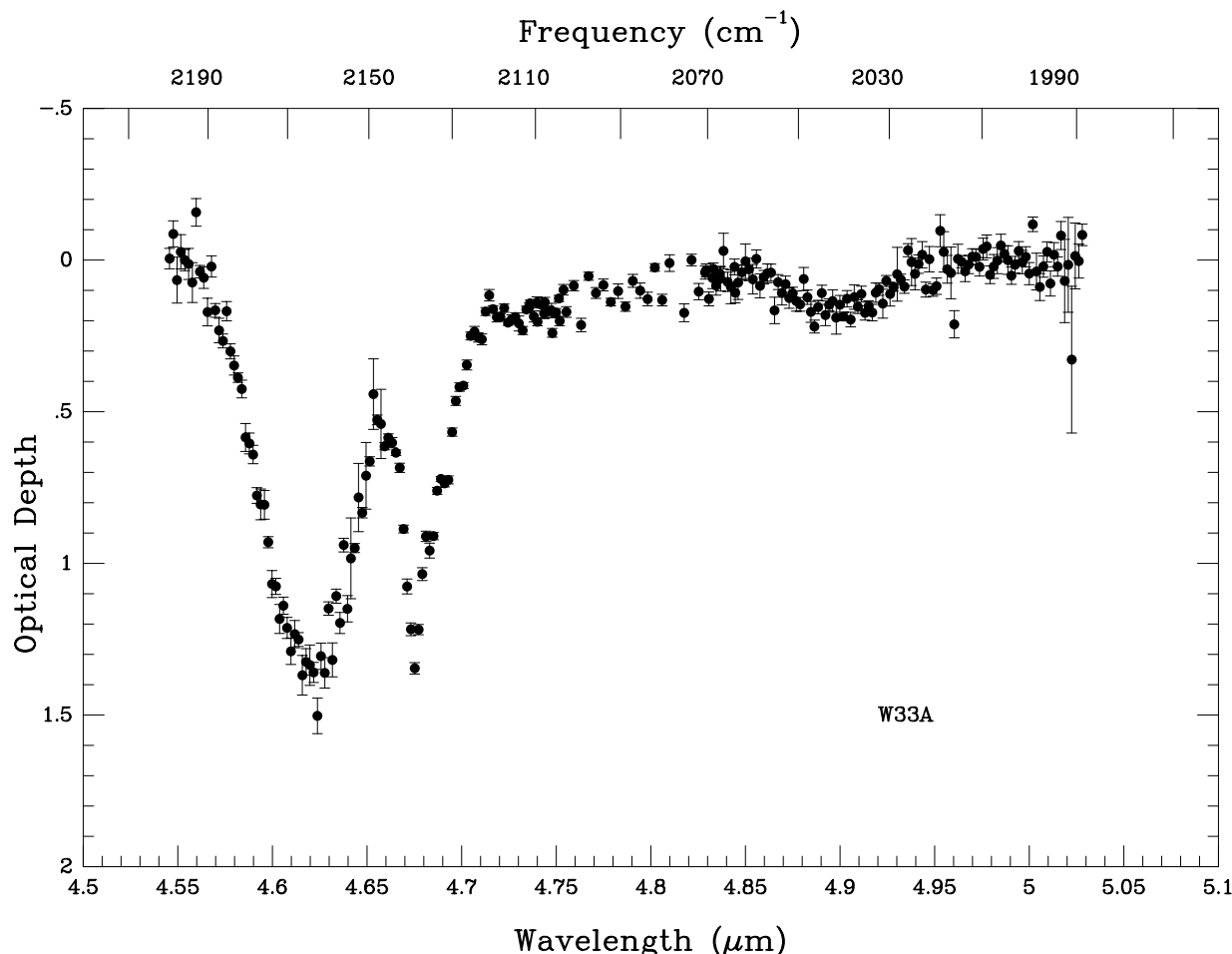


FIG. 4.—Optical depth plot from 4.5–5.1  $\mu\text{m}$  (2222–1960  $\text{cm}^{-1}$ ) for W33 A

thiocyanates ( $\text{X}-\text{N}\equiv\text{C}-\text{S}$ ), and cyanate ions ( $\text{X}-\text{O}-\text{C}\equiv\text{N}$ ) are among those that appear in the 4–5  $\mu\text{m}$  (2500–2000  $\text{cm}^{-1}$ ) region. Cumulative double bond species are also active in this region, and these include isocyanate ions ( $\text{X}-\text{N}=\text{C}=\text{O}$ ) and ketenes ( $\text{C}=\text{C}=\text{O}$ ). Cyanate ions will be discussed separately from the other candidates in the following section, as there has been substantially more laboratory investigation of these particular candidates.

#### 4.1. Nitriles and Isonitriles

Most nitriles absorb at much higher frequencies (2200–2280  $\text{cm}^{-1}$ ) than the interstellar 4.62  $\mu\text{m}$  (2165  $\text{cm}^{-1}$ ) band and are therefore not good candidate carriers (d’Hendecourt et al. 1986). In these cases, the spectra are not very sensitive to the matrix environment, so that there is little expectation that the peak wavelength might be shifted to the interstellar frequency in a different ice “environment.” Among the specific nitriles that can be ruled out as the identification of the interstellar 4.62  $\mu\text{m}$  band are dicyandiamide [ $(\text{NHC}(\text{NH}_2)\text{NHCN})$ ] and cyanogen ( $\text{C}_2\text{N}_2$ ), because they have longer wavelength features that are not present in the observational data. Pyruvonitrile ( $\text{CH}_3\text{COCN}$ ) had been proposed (d’Hendecourt et al. 1986), but subsequent experiments make this assignment seem unlikely (Bernstein et al. 1997). Nitriles that form strong hydrogen bonds with  $\text{H}_2\text{O}$  are an exception, however, because a substantial shift in the peak absorption frequency is seen in these cases. In the  $\text{X}-\text{C}\equiv\text{N}$  configu-

ration, the presence of a lone electron pair on X allows hydrogen to bond with that atom, adjacent to the nitrile, and because of electron sharing between the nitrile and the X group, the H bond affects the peak position of the nitrile band (Bernstein et al. 1997).

Isonitriles absorb in the 2150–2175  $\text{cm}^{-1}$  spectral range and, as such, are good candidates for the carrier of the interstellar 2165  $\text{cm}^{-1}$  band (d’Hendecourt et al. 1986; Larson et al. 1985). Methyl isocyanide ( $\text{CH}_3\text{NC}$ ; Larson et al. 1985) is such a candidate, because both in the gas phase and in an argon matrix  $\text{CH}_3\text{NC}$  absorbs at 2161  $\text{cm}^{-1}$  (Khlifi et al. 1996; Freedman & Nixon 1972). The isonitriles, *n*-butylisocyanide [ $\text{CH}_3(\text{CH}_2)_3\text{NC}$ ] and methylisocynoacetate ( $\text{CH}_3\text{OCOCH}_2\text{NC}$ ), have also been considered as carriers of the interstellar 2165  $\text{cm}^{-1}$  band (Bernstein et al. 1997). Both have a strong CN stretch near the interstellar 2165  $\text{cm}^{-1}$  band in argon and  $\text{H}_2\text{O}$  matrices (Bernstein et al. 1997). However, one general, indirect argument against all isocyanides is that their abundance might reasonably be expected to be less than that of the corresponding cyanide, and no bands in the general nitrile spectral region (where cyanides would appear) have been seen (Whittet et al. 1996). For this reason, the feasibility label for isonitriles in Table 4 is “weak.”

#### 4.2. Species with Cumulative Double Bonds

Among species with cumulative double bonds are ketenes ( $\text{C}=\text{C}=\text{O}$ ). Ketenes absorb in the range 2080–2200  $\text{cm}^{-1}$ ,

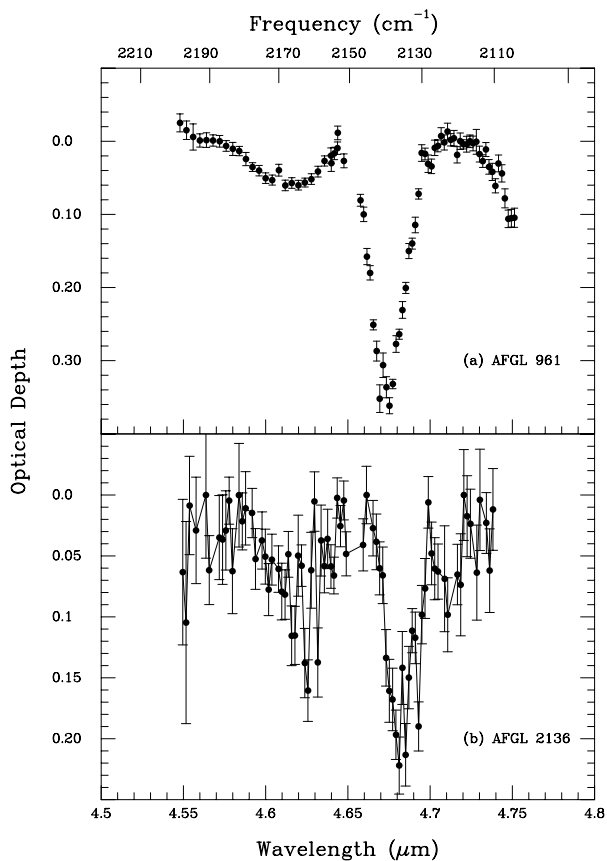


FIG. 5.—Optical depth plot from 4.5–4.8  $\mu\text{m}$  (2222–2083  $\text{cm}^{-1}$ ) for (a) AFGL 961 and (b) AFGL 2136.

often near 4.65  $\mu\text{m}$  (2150  $\text{cm}^{-1}$ ) (Socrates 1980). In an argon or  $\text{O}_2$  matrix, the compound ketene ( $\text{H}_2\text{C}_2\text{O}$ ) absorbs at 4.67  $\mu\text{m}$  (2142  $\text{cm}^{-1}$ ) (Moore & Pimentel 1963), which is close enough to the interstellar position to speculate that in

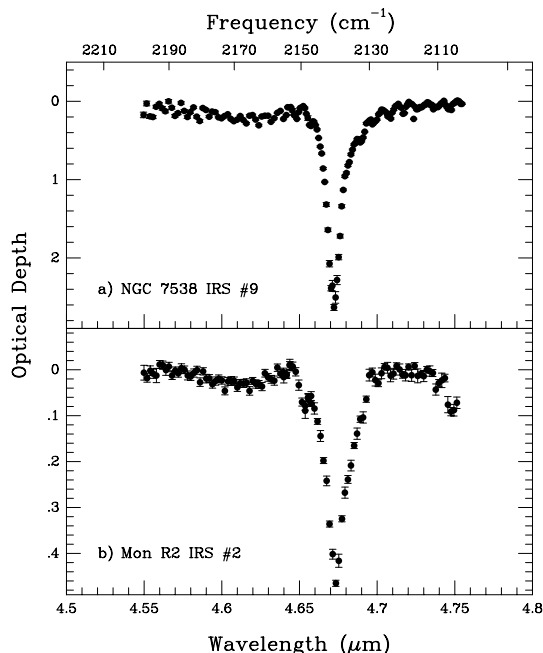


FIG. 6.—Optical depth plot from 4.5–4.8  $\mu\text{m}$  (2222–2083  $\text{cm}^{-1}$ ) for (a) NGC 7538 IRS 9 and (b) Mon R2 IRS 2.

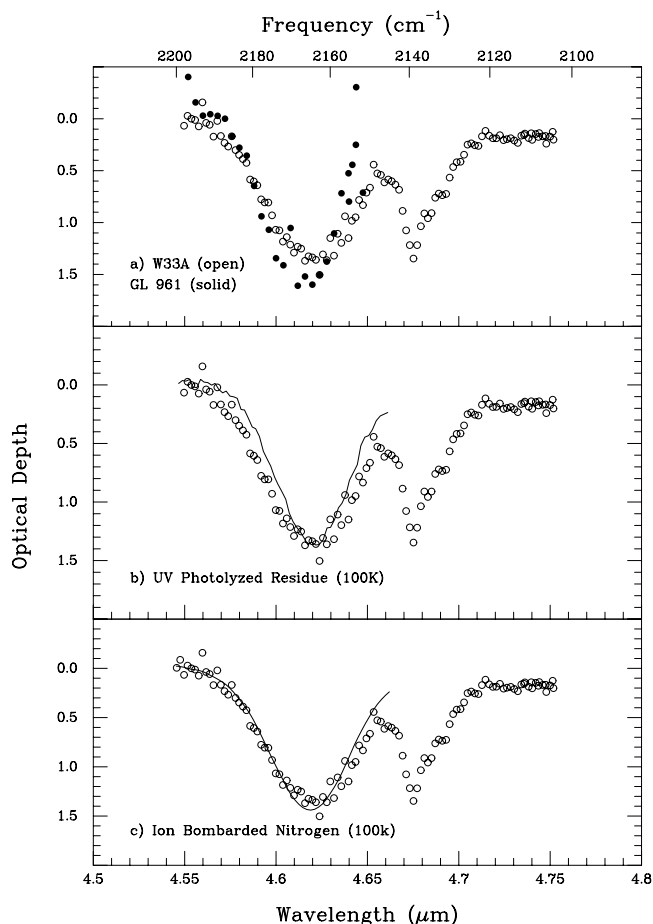


FIG. 7.—The 4.5–4.8  $\mu\text{m}$  (2222–2083  $\text{cm}^{-1}$ ) flux spectra of (a) AFGL 961 (*filled circles*) scaled to W33 A (*open circles*); for clarity, only the points relevant to the 4.62  $\mu\text{m}$  band are shown for W33 A. (b) The UV photolyzed residue produced when an  $\text{H}_2\text{O}:\text{CH}_3\text{OH}:\text{CO}:\text{NH}_3 = 100:50:10:10$  ice is irradiated at 12 K and subsequently warmed to 100 K (taken from Bernstein et al. 1995; *solid line*) compared to the observational data from W33 A (*open circles*). (c) The organic residue produced through 3 keV  $\text{He}^+$  ion bombardment of a mixture containing  $\text{H}_2\text{O}:\text{N}_2:\text{CH}_4 = 1:1:1$  compared to the observational data for W33 A (*open circles*). The dose received by the mixture was 40 eV per 16 amu, and the spectrum shown is after warm-up to 100 K (Palumbo et al. 1998) (*solid line*).

an  $\text{H}_2\text{O}$  matrix ketene might be a viable candidate. Laboratory experiments to investigate the peak position of ketene absorption in astrophysically relevant matrices (such as  $\text{H}_2\text{O}$ ,  $\text{CO}$ , and  $\text{CO}_2$ ) are needed.

### 4.3. The Cyanate Ion

The strong CN stretching vibration in simple cyanate salts falls in the range 4.5–4.62  $\mu\text{m}$  (2225–2165  $\text{cm}^{-1}$ ). For example,  $\text{NH}_4\text{NCO}$  absorbs at 2217  $\text{cm}^{-1}$  (Hubbard et al. 1975) Grim & Greenberg (1987a); Schutte & Greenberg (1997); Demyk et al. (1999) have identified the 4.62  $\mu\text{m}$  band in their laboratory spectra, generated by FUV photolysis, as the cyanate ( $\text{OCN}^-$ ) ion and suggest that this ion is the carrier of the interstellar band. The laboratory spectra compare favorably with the interstellar 4.62  $\mu\text{m}$  band (Fig. 7b), and for this reason we have assigned its feasibility as “good” in Table 4.

### 4.4. Thiocyanates

The  $\text{C}\equiv\text{N}$  stretching vibrations in thiocyanates ( $\text{S}-\text{TC}\equiv\text{N}$ ) is observed in the range 2140–2175  $\text{cm}^{-1}$



(Socrates 1980). There have been no studies made thus far to test their relevance to the interstellar  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) band.

#### 4.5. Silanes

Organic silane ( $-\text{SiH}_3$ ) compounds have also been considered candidates for the interstellar  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) band (Nuth & Moore 1988; Moore, Tanabe, & Nuth 1991). The Si–H stretch in organic silanes ( $-\text{SiH}_3$ ) falls in the range  $2160\text{--}2140 \text{ cm}^{-1}$  (Socrates 1980). Silane ( $\text{SiH}_4$ ) has two absorption bands, at  $2216$  and  $2176 \text{ cm}^{-1}$  (in argon) and  $2200$  and  $2165 \text{ cm}^{-1}$  (in  $\text{H}_2\text{O}$ ) (Lloret & Abouaf-Marguin 1986; Nuth & Moore 1988). Only one band is detected in the interstellar spectra, and therefore  $\text{SiH}_4$  itself is not a good candidate. In addition, we note that hydrogen bonded to silicon surfaces shows Si–H stretching modes in the range  $2200\text{--}2000 \text{ cm}^{-1}$ ; the exact position will depend on the molecular group attached. For silicate surfaces, the oxidized monohydride  $[\text{Si}_3\text{O}_3\text{Si}-\text{H}]$  absorbs at  $2250 \text{ cm}^{-1}$ . For silicon surfaces, the monohydride  $[\text{Si}_3\text{Si}-\text{H}]$  absorbs at  $2082 \text{ cm}^{-1}$ , and the dihydride  $[\text{Si}_2\text{SiH}_2]$  absorbs at  $2133 \text{ cm}^{-1}$  (Watanabe & Sugati 1996). Presumably there is a silicon hydride that falls near  $2165 \text{ cm}^{-1}$ , but until laboratory spectra reveal such a band, the feasibility of silanes must be labeled “weak” in Table 4.

#### 4.6. Summary of Candidate Species

In the literature, a number of likely candidates have been proposed for the interstellar  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) band (see Table 4). A few of the specific proposed carriers can be confidently eliminated (dicyandiamide, cyanogen, and  $\text{SiH}_4$ ) because the stronger, corresponding infrared absorption bands of these species have not been observed. Others, such as acetonitrile and pyruvitrile, are rejected because they absorb at the wrong wavelengths. A number of other general types of compounds are candidates (i.e., isonitriles and organic silanes), but a likely specific molecule has not been put forward and further study is required. Additional candidates that require laboratory study include thiocyanates, isothiocyanates, and ketenes. The cyanate ion may be the most promising candidate of all. It is always difficult to identify with confidence a compound from one spectral feature, and hence a true identification of the carrier of the  $4.62 \mu\text{m}$  band still cannot be made on spectroscopic arguments alone. Other features have been searched for, but within the limited data available (Tables 2 and 3) no other absorption feature correlates with the interstellar  $4.62 \mu\text{m}$  band. As the data reduction techniques for evaluation of the data from *ISO* improve, it is likely that weaker features will be probed more effectively, and bands related to the  $2165 \text{ cm}^{-1}$  carrier may be identified. This will greatly assist in the assessment of the possible carriers presently under consideration.

### 5. PRODUCTION PATHWAYS TO THE INTERSTELLAR 4.62 MICRON BAND

The pathway toward production of the interstellar  $4.62 \mu\text{m}$  band is a complex issue. From the laboratory data currently available,  $4.62 \mu\text{m}$  bands similar in profile can be produced from ice mixtures containing nitrogen, carbon, and oxygen, which are exposed to either ultraviolet photolysis or ion bombardment. Accepting, for the moment, the identification of the interstellar  $4.62 \mu\text{m}$  band with the

$4.62 \mu\text{m}$  band produced by processing ice analogs, we will now discuss the possible pathways toward the interstellar feature. Energetic processing of interstellar ices, either through ultraviolet photolysis, ion bombardment, or both, and grain surface chemistry are all plausible pathways by which the interstellar  $4.62 \mu\text{m}$  band might arise. Only ion bombardment, however, has been shown to form the  $4.62 \mu\text{m}$  band if the nitrogen is in the form of solid  $\text{N}_2$ . The discovery that solid  $\text{N}_2$  comprises the majority of the ice present on bodies in the outer solar system (Cruikshank et al. 1993, 1998) suggests that the infrared inactive  $\text{N}_2$  may be hiding in the interstellar ices as well (Ehrenfreund et al. 1997; Elsila, Allamandola, & Sandford 1997). Therefore, studies of the interstellar  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) band may provide information on both the reservoir of the solid-state nitrogen in interstellar ices and the energetic process by which it participates in the chemistry of the carrier of the  $4.62 \mu\text{m}$  band. Identifying the production mechanism is therefore directly related to identifications of the interstellar ice constituents that reside on the preprocessed grain mantle. Here we discuss each of these in turn, and they are summarized in Table 5.

#### 5.1. FUV Laboratory Studies

The role of ultraviolet photolysis is one of the key questions in the chemistry of interstellar ices. Theoretical and experimental studies have long suggested that FUV photolysis of interstellar ices is important for their composition and the X–CN band is often cited in support of this hypothesis (Greenberg & Yenchu 1973; d’Hendecourt et al. 1986; Allamandola, Sandford, & Valero 1988; Grim et al. 1989a; Grim & Greenberg 1987; Schutte et al. 1996; Bernstein et al. 1996). Laboratory far-ultraviolet (FUV) photolysis of ices containing CO and  $\text{NH}_3$  show a prominent absorption feature at  $4.62 \mu\text{m}$  ( $2165 \text{ cm}^{-1}$ ) (Hagen 1982; Lacy et al. 1984; d’Hendecourt et al. 1986; Grim 1988; Grim & Greenberg 1987a; Allamandola et al. 1988; Schutte & Greenberg 1997). Figure 7b illustrates the good comparison between the FUV photolysis laboratory and interstellar data.

While the first experiments were interpreted in terms of an otherwise unidentified nitrile, very soon afterward the ion  $\text{OCN}^-$  was implicated (Grim & Greenberg 1987a; Grim et al. 1989a; Schutte & Greenberg 1997). Recent laboratory studies using electron donors and acceptors support this ion identification for the laboratory-produced  $4.62 \mu\text{m}$  band (Demyk et al. 1999). The ion is thought to result from an acid-base reaction between  $\text{NH}_3$  and FUV-produced HNC(O) (Grim et al. 1989a; Schutte & Greenberg 1997).

One problem for any ion identification is the charge balance. For each negative ion made, a positive ion must be produced. In these experiments, the proposed counter-ion is  $\text{NH}_4^+$ , and a band in the range  $6.6\text{--}6.9 \mu\text{m}$  ( $1450\text{--}1500 \text{ cm}^{-1}$ ), depending on the matrix, is attributed to this ion (Grim & Greenberg 1987a; Demyk et al. 1999). There is a strong band in interstellar ice spectra at about  $6.8 \mu\text{m}$ , but whether the  $\text{NH}_4^+$  band falls at this position in astrophysically relevant matrices is not clear yet. Nor has the charge balance in interstellar ices been completely evaluated. The sensitivity of the strength and peak position of the  $\text{NH}_4^+$  band to the molecular environment in the laboratory studies (Demyk et al. 1999) might indicate that such an analysis will be difficult.

### 5.2. Ion Bombardment of Ices

An alternative way to process interstellar ice mixtures is through bombardment with high-energy protons. Moore et al. (1983) have shown that a band is produced at  $2170\text{ cm}^{-1}$  upon the irradiation of a mixture of  $\text{H}_2\text{O}:\text{NH}_3:\text{CH}_4$  ( $\approx 1:3:2$ ) after irradiation with  $\approx \text{MeV}$  protons. They also found that irradiation of mixtures  $\text{H}_2\text{O}:\text{N}_2:\text{CO}_2$  ( $\approx 1:1:1$ ) and  $\text{H}_2\text{O}:\text{N}_2:\text{CO}$  ( $\approx 5:1:1$ ) produces a band at about  $2180\text{ cm}^{-1}$ . It is well established (e.g., Strazzulla & Johnson 1991; Strazzulla 1997, 1998) that ion irradiation drives the formation not only of specific molecules, depending on the particular target, but also of refractory residues. Especially promising is recent work (e.g., Palumbo et al. 1998) that includes  $\text{N}_2$  in the starting mixture and produces a band that is quite similar to the interstellar band. Ion bombardment, which is analogous to cosmic-ray interactions with icy dust grains in dense molecular clouds, has the advantage that the high-energy ions can break  $\text{N}_2$ . The resulting  $4.62\text{ }\mu\text{m}$  band is discussed in Palumbo et al. (1998) and compares well to the interstellar data for W33 A (Fig. 7c). It is unclear whether the  $4.62\text{ }\mu\text{m}$  band produced by ion bombardment is also due to the  $\text{OCN}^-$  ion.

Laboratory studies have shown that UV photolysis of  $\text{N}_2$ -containing mixtures cannot produce a  $4.62\text{ }\mu\text{m}$  band (Elsila et al. 1997), because the energy from the UV lamps is insufficient to break apart the  $\text{N}_2$ . While the interstellar radiation field contains photons more energetic (i.e., in the range  $912\text{--}1250\text{ \AA}$ ) than these laboratory lamps, FUV photolysis at best will be very inefficient.

### 5.3. Grain Surface Reactions

It is also possible that the carrier of the  $4.62\text{ }\mu\text{m}$  band is produced via simple grain surface chemistry, followed by the acid-base reactions outlined in § 5.1 (Keane & Schutte 1999). Laboratory studies have shown that even in the absence of FUV irradiation codeposited HNCO and  $\text{NH}_3$ , in various astrophysical relevant ice mixtures, undergoes an acid-base reaction forming  $\text{OCN}^-$  and  $\text{NH}_4^+$  (Keane 1997; Keane & Schutte 1999). While not all of the HNCO is converted in these experiments, upon warm-up to  $80\text{ K}$  much of the remaining HNCO undergoes the acid-base reaction, and the spectrum is dominated by the  $4.62\text{ }\mu\text{m}$  band. HNCO and  $\text{NH}_3$  might both result from grain surface chemistry. Ammonia can be formed through hydrogenation of accreted atomic N. HNCO might result from chemistry initiated by the reaction of  $\text{H} + \text{CO}$  forming HCO (van IJzendoorn et al. 1983). This radical can react with N and the intermediate excited complex  $\text{NCHO}^*$  might relax to the energetically most favorable isomer HNCO (Charnley 1997; 1999). Alternatively, this route might actually lead to  $\text{NH}_2\text{CHO}$  (Tielens & Hagen 1982), a species that has been observed in the gas phase in the ISM (Kuan & Snyder 1996), but which has no solid-state  $4.62\text{ }\mu\text{m}$  band. Further studies of this grain surface route are needed.

## 6. CONCLUSIONS

We present new infrared spectra of dust embedded protostars that reveal a strong absorption feature near  $4.67\text{ }\mu\text{m}$  ( $2141\text{ cm}^{-1}$ ) and a weaker absorption feature centered near  $4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ). The feature near  $4.67\text{ }\mu\text{m}$  ( $2141\text{ cm}^{-1}$ ) has previously been identified as absorption due to solid CO within the icy grain mantle (Tielens et al. 1991). The

$4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ) band, however, has not yet been unambiguously identified. The band present throughout our sample of five embedded protostars was previously reported in a small number of other protostellar objects (Lacy et al. 1984; Tegler et al. 1993; Tegler et al. 1995; Weintraub et al. 1994), demonstrating that the feature probably forms fairly readily in protostellar regions. The interstellar  $4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ) band does not vary more than  $2.5\text{ cm}^{-1}$  in position, suggesting the carrier of the band may be fairly insensitive to changes in the local environment. Gas-phase CO absorption is detected in the spectra of some of these sources, but blending of gas-phase CO is not responsible for the broad, shallow absorption at  $4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ). The  $4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ) band does not correlate with any other absorption feature in the spectra of these sources; however, there is in general a lack of correlation among the ice absorption features in these sources. The  $4.62\text{ }\mu\text{m}$  band may therefore simply be a part of that same trend, perhaps indicating the carriers of these features are sensitive to environmental conditions in the cloud.

Our study samples protostars from different molecular clouds and shows that the strength of the  $4.62\text{ }\mu\text{m}$  band per unit  $A_V$  (total visual extinction) varies by about an order of magnitude from cloud to cloud. Such a disparity may indicate that different processes (i.e., UV photolysis, ion bombardment, grain surface chemistry, or thermal processes) are involved in the production of the  $4.62\text{ }\mu\text{m}$  band or, more likely, that differing evolutionary states of the embedded objects obfuscate the situation. It is noteworthy that in every case in which the  $4.62\text{ }\mu\text{m}$  feature was searched for, an absorption band was found.

We have investigated a variety of possible identifications for the interstellar  $4.62\text{ }\mu\text{m}$  band. Good fits to the interstellar spectra are obtained with the  $\text{OCN}^-$  ion produced either by acid-base reactions (Keane & Schutte 1999) or through FUV photolysis of  $\text{NH}_3$ -containing ices (Schutte & Greenberg 1997; Demyk et al. 1999), as well as with an organic residue produced through ion bombardment of nitrogen-containing ices (Moore et al. 1983; Palumbo et al. 1998). A number of other species may be feasible candidates as well, but relevant laboratory data are incomplete.

We have reviewed proposed formation mechanisms for the production of the  $4.62\text{ }\mu\text{m}$  band, such as UV photolysis, ion bombardment, and/or grain surface chemistry assisted by acid-base reactions in interstellar ices. Organic residues produced by energetic processing of astrophysically plausible initial ice mixtures reproduce the  $4.62\text{ }\mu\text{m}$  ( $2165\text{ cm}^{-1}$ ) band in general peak position and profile quite well. Ion bombardment has the advantage that it would allow  $\text{N}_2$ , possibly the dominant nitrogen reservoir in molecular clouds, to participate in the chemistry of the ices. The issues to consider in trying to discriminate between these processes are the nitrogen budget in interstellar ices (i.e.,  $\text{NH}_3$  or  $\text{N}_2$ ), the efficiency of FUV photolysis versus particle bombardment, and the detailed grain surface reaction schemes. None of these are well determined at this time.

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